

3rd YEAR CMP COMPULSORY EXPERIMENT

Introduction to the Ising model

TASK:

1. Show that the lowest possible energy for the Ising model is $E = -DNJ$, where D is the number of dimensions and N is the total number of spins.
2. What is the multiplicity of this state?
3. Calculate its entropy.

1.

$$E = -\frac{1}{2}J \sum_i^N \sum_{j \in \text{neighbours}(i)} s_i s_j$$

In the ground state all spins face the same direction (all +1 or all -1). By trail:

Dimension	$\sum_{j \in \text{neighbours}(i)} s_i s_j$ in ground state
1D	2
2D	4
3D	6

Gives the relationship:

$$\sum_{j \in \text{neighbours}(i)} s_i s_j = 2D$$

$$\sum_i^N \sum_{j \in \text{neighbours}(i)} s_i s_j = N(2D)$$

Therefore:

$$E = -\frac{1}{2}J \sum_i^N \sum_{j \in \text{neighbours}(i)} s_i s_j = -\frac{1}{2}J \times N(2D)$$

And so the lowest possible energy is:

$$E = -DNJ$$

2. The multiplicity of this state is 2.
3. Definition of entropy:

$$S = k_B \ln W$$

Where W is number of microstates, k_B is the Boltzmann constant and S is the entropy.

Since $W = 2$:

$$S^0 = 1.38 \times 10^{-23} \times \ln 2$$

$$S^0 = 9.565 \times 10^{-24} \text{ J K}^{-1}$$

TASK: Imagine that the system is in the lowest energy configuration. To move to a different state, one of the spins must spontaneously change direction ("flip").

1. What is the change in energy if this happens ($D = 3$, $N = 1000$)?
2. How much entropy does the system gain by doing so?

1. Lowest energy configuration is all spins aligned, where $E = -DNJ$

1 spin flipped:

$$E = -\frac{1}{2}J((2ND) - 6 - 6)$$

$$E = -J(ND - 6)$$

because loss of favourable interaction gives +6, and gain of unfavourable interaction gives +6 again since $\downarrow\uparrow$ is unfavourable. The number 6 is due to each cell having 6 adjacent cells in a 3-D structure and so each $6 * -1 = -6$ as the cell has opposite spin to the adjacent cells.

When $D = 3$ and $N = 1000$:

$$E = -J * (3 * 1000 - 6)$$

$$E = -2994J$$

Change in E between ground state and spin flipped state:

$$\Delta E = E - E^0$$

$$\Delta E = -2994J - -3000J = +6J$$

2. The entropy becomes:

$$W = N = 1000$$

$$S = 1.38E - 23 * \ln 1000 = 9.533E - 23 \text{ JK}^{-1}$$

Therefore the change in entropy is:

$$\Delta S = S - S^0$$

$$\Delta S = 9.533E - 23 - 9.565E - 24 = 8.576E - 23 \text{ JK}^{-1}$$

$$\Delta S = 51.63 \text{ JK}^{-1}\text{mol}^{-1}$$

Hence the system has increased in entropy as it has become more disordered.

TASK:

1. Calculate the magnetisation of the 1D and 2D lattices in figure 1.
2. What magnetisation would you expect to observe for an Ising lattice with $D = 3$, $N = 1000$ at absolute zero?

1. Expression for magnetization is given by:

$$M = \sum_i s_i$$

$$1D \text{ lattice: } M = 3 - 2 = +1$$

$$2D \text{ lattice: } M = 13 - 12 = +1$$

2. At 0 K the system will be in its ground state. Therefore all the spins will be in the same directions (all +1 or all -1).

$$D = 3, N = 1000$$

$$M = \pm 1000$$

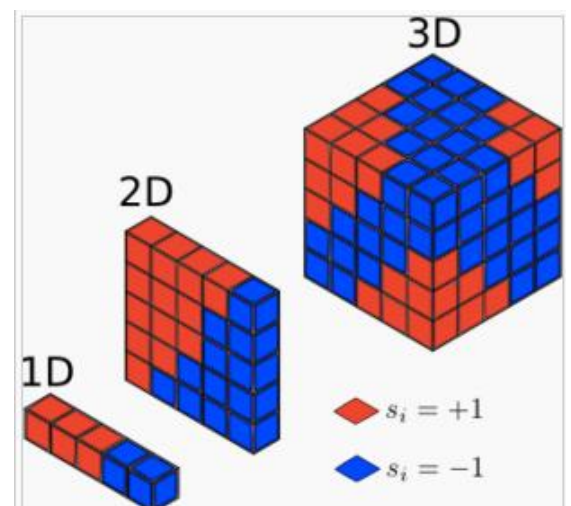


Figure 1. Illustration of an Ising Lattice in 1D, 2D and 3D. Red cells are +1 and blue cells are -1

Calculating the energy and magnetisation

TASK: complete the functions `energy()` and `magnetisation()`, which should return the energy of the lattice and the total magnetisation, respectively. In the `energy()` function you may assume that $J = 1.0$ at all times (in fact, we are working in *reduced units* in which $J = k_B$, but there will be more information about this in later sections). Do not worry about the efficiency of the code at the moment — we will address the speed in a later part of the experiment.

The code can be found here:

[EnergyAndMagnetisation.py](#)

TASK: Run the `ILcheck.py` script from the IPython Qt console using the command

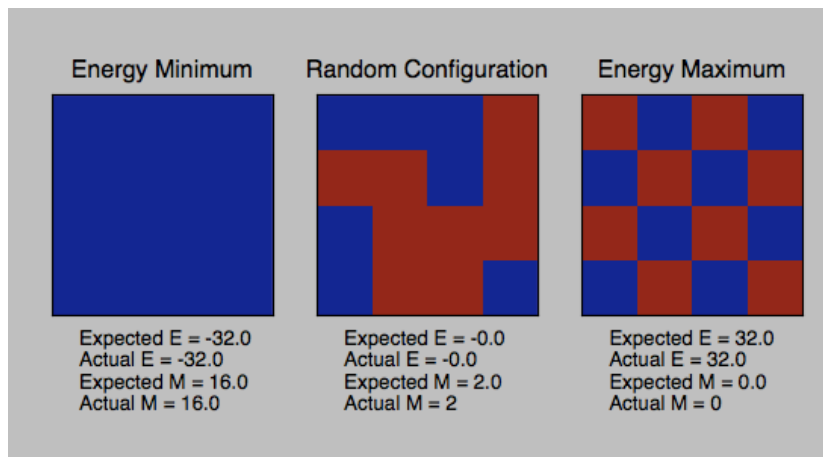


Figure 2. output of `ILcheck.py`. Shows the script is successful

Introduction to Monte Carlo simulation

TASK:

1. How many configurations are available to a system with 100 spins?
2. To evaluate these expressions, we have to calculate the energy and magnetisation for each of these configurations, then perform the sum. Let's be very, very, generous, and say that we can analyse 1×10^9 configurations per second with our computer. How long will it take to evaluate a single value of $\langle M \rangle_T$?

1. Since there are 2 options of what each cell could be (+1 or -1):

$$\text{no. of possible configs} = 2^{100} = 1.26 \times 10^{30}$$

Time to evaluate configurations of 100 spin system, T:

$$T = \frac{1.26 \times 10^{30}}{1 \times 10^9} = 1.268 \times 10^{21} \text{ seconds}$$

$$T = 40,181,353,532,343 \text{ years}$$

TASK: Implement a single cycle of the above algorithm in the `montecarlocycle(T)` function. This function should return the energy of your lattice and the magnetisation at the end of the cycle. You may assume that the energy returned by your `energy()` function is in units of k_B ! Complete the `statistics()` function.

This should return the following quantities whenever it is called: $\langle E \rangle$, $\langle E^2 \rangle$, $\langle M \rangle$, $\langle M^2 \rangle$, and the number of Monte Carlo steps that have elapsed.

The code can be found here:

[MonteCarloAndStatistics.py](#)

TASK:

1. If $T < T_C$, do you expect a spontaneous magnetisation (i.e. do you expect $\langle M \rangle \neq 0$)?
2. When the state of the simulation appears to stop changing (when you have reached an equilibrium state), use the controls to export the output to PNG and attach this to your report.
3. You should also include the output from your statistics() function.

1. The Curie temperature is the temperature at which permanent magnetism changes to unduced magnetism. If Temp is less than the Curie temperature then expect $\langle M \rangle \neq 0$ as all the spins will be aligned in a ferromagnetic material so there will be a strong dipole. 8x8 lattice run in figure below:

2.

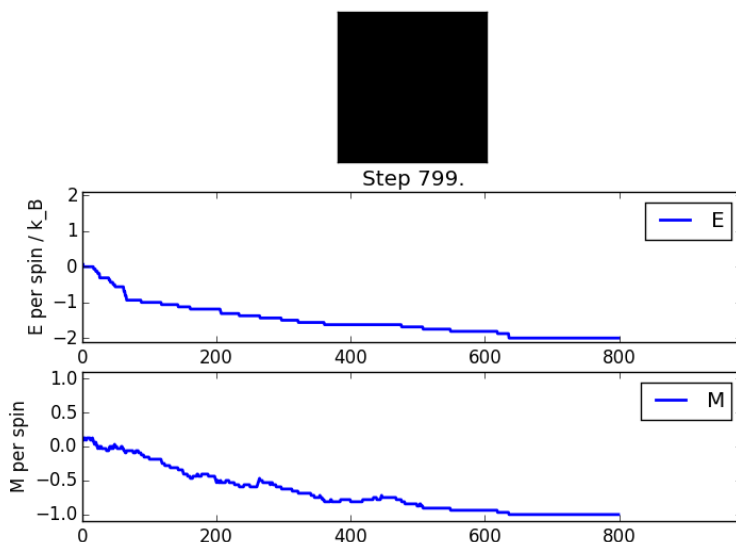


Figure 3. Output of ILanim after reaching equilibrium. All cells are aligned (all one colour)

3.

```
In [9]: %run ILanim.py
Averaged quantities:
E = -1.43842645382
E*E = 2.29835376283
M = -0.667830672748
M*M = 0.532800028506
```

Figure 4. Output values of ILanim.py. showing, from top: Energy, energy squared, magnetisation and magnetisation squared

Accelerating the code

TASK: Use the script ILtimetrial.py to record how long your *current* version of IsingLattice.py takes to perform 2000 Monte Carlo steps. This will vary, depending on what else the computer happens to be doing, so perform repeats and report the error in your average!

Time to perform 2000 steps/ seconds
8.111257549
8.084071129
8.090109769
8.047899733
8.036258632
8.047862672
8.083449932
Standard Deviation: 0.02764491
Standard Error: 0.010448794
Average: 8.071558488 ± 0.010448794

TASK: Look at the documentation for the NumPy sum function. You should be able to modify your magnetisation() function so that it uses this to evaluate M. The energy is a little trickier. Familiarise yourself with the NumPy roll and multiply functions, and use these to replace your energy double loop (you will need to call roll and multiply twice!).

The code can be found here:

[AcceleratedEnergyAndMagnetisation.py:](#)

TASK: Use the script ILtimetrial.py to record how long your *new* version of IsingLattice.py takes to perform 2000 Monte Carlo steps. This will vary, depending on what else the computer happens to be doing, so perform repeats and report the error in your average!

Accelerated Code
0.182258169
0.190376778
0.179950703
0.179642813
0.180751218
0.208782055
0.17884857
Standard Deviation: 0.010859743
Standard Error: 0.004104597
Average: 0.185801472 ± 0.004104597

The code has dramatically increased in efficiency.

The effect of temperature

TASK: The script `ILfinalframe.py` runs for a given number of cycles at a given temperature, then plots a depiction of the *final* lattice state as well as graphs of the energy and magnetisation as a function of cycle number. This is much quicker than animating every frame!

1. Experiment with different temperature and lattice sizes. How many cycles are typically needed for the system to go from its random starting position to the equilibrium state?
2. Modify your `statistics()` and `montecarlostep()` functions so that the first N cycles of the simulation are ignored when calculating the averages.
3. You should state in your report what period you chose to ignore, and include graphs from `ILfinalframe.py` to illustrate your motivation in choosing this figure.

1.

Steps to reach equilibrium
450
520
800
2800
400
890
650

Below is the first section of the `ILfinalframe.py` expanded to find where it reaches equilibrium (where the line is flat). In this case the value is 650. This technique was repeated, giving the table above.

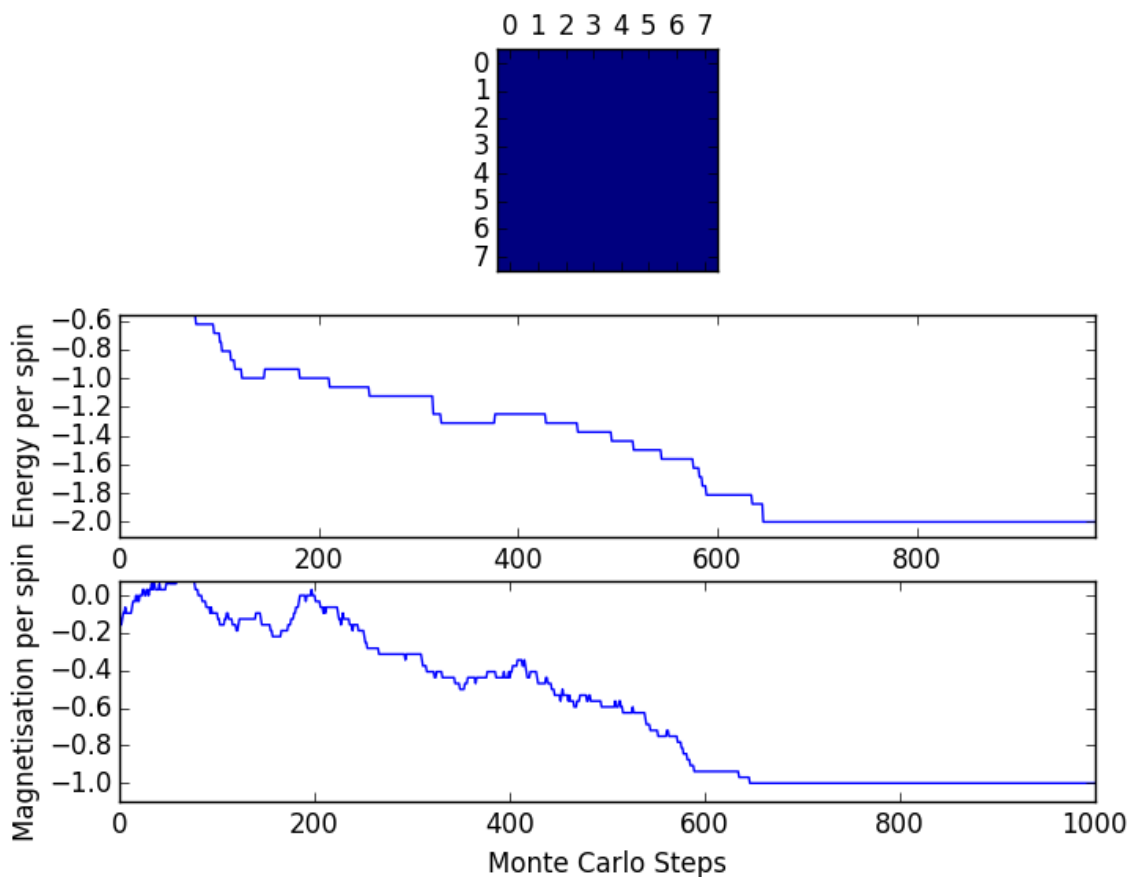


Figure 5. Expanded view of `ILfinalframe.py` to find what no. cycle equilibrium is reached for 8x8 lattice.

2. The code can be found here:

[ModifiedMonteCarloAndStastics.py](#)

3. Maximum was 2800 so 4000 was chosen as the cut off as it is significantly bigger than 2800, however it isn't so big so that it takes a lot of time to get results. This is not very accurate as only an 8x8 lattice was investigated.

TASK: Use `ILtemperaturerange.py` to plot the average energy and magnetisation for each temperature, *with error bars*, for an 8×8 lattice. Use your intuition and results from the script `ILfinalframe.py` to estimate how many cycles each simulation should be. The temperature range 0.25 to 5.0 is sufficient. Use as many temperature points as you feel necessary to illustrate the trend, but do not use a temperature spacing larger than 0.5. The NumPy function `saveetxt()` stores your array of output data on disk — you will need it later. Save the file as `8x8.dat` so that you know which lattice size it came from.

The code can be found here:

[ILtemperaturerange.py](#)

The number of cycles I chose is 200,000 as that means that the cut off (4000) is only 2% of the total cycles. which is a small proportion. There are small fluctuations which will have little effect on the average as they are such a small proportion of the total.

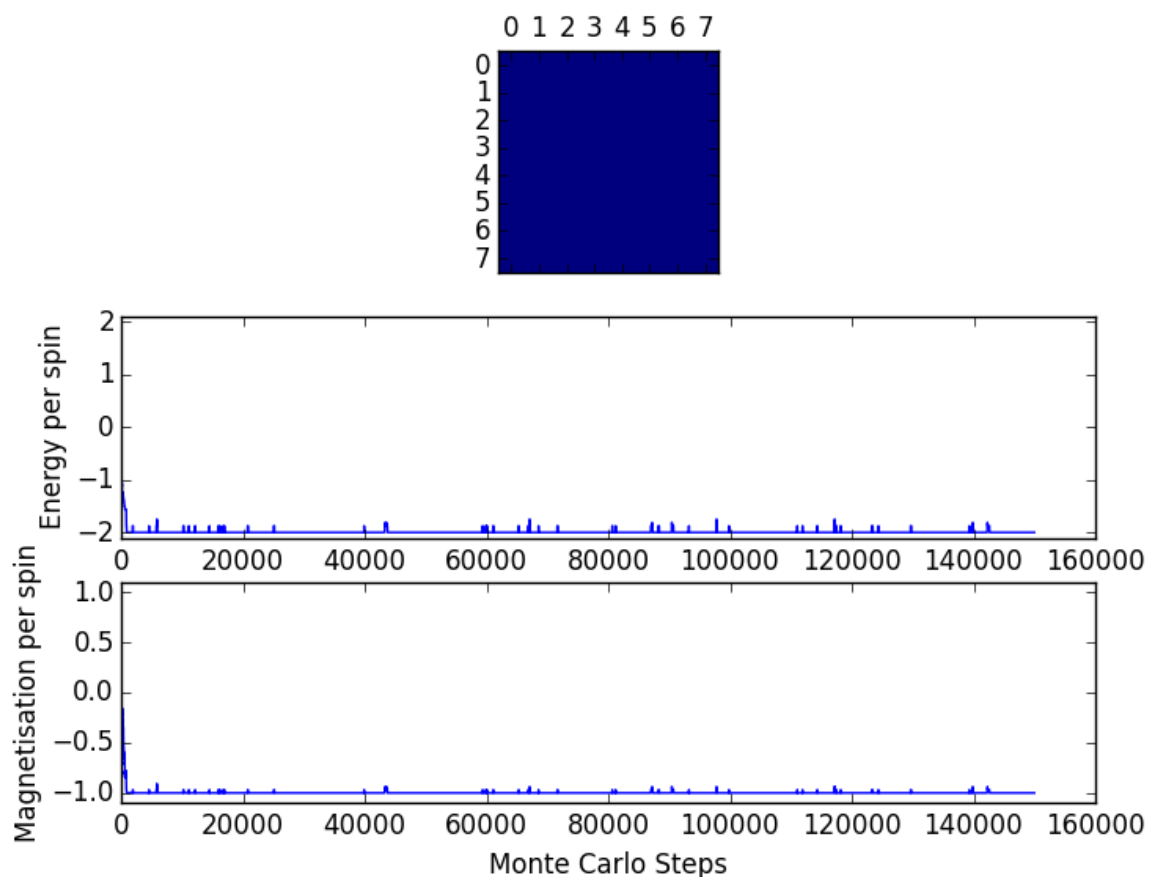


Figure 6. *Non Expanded output of ILfinalframe.py for 8x8 lattice*

The plot below shows the energy per spin vs temperature for the 8x8 lattice with error bars. The error bars are very small.

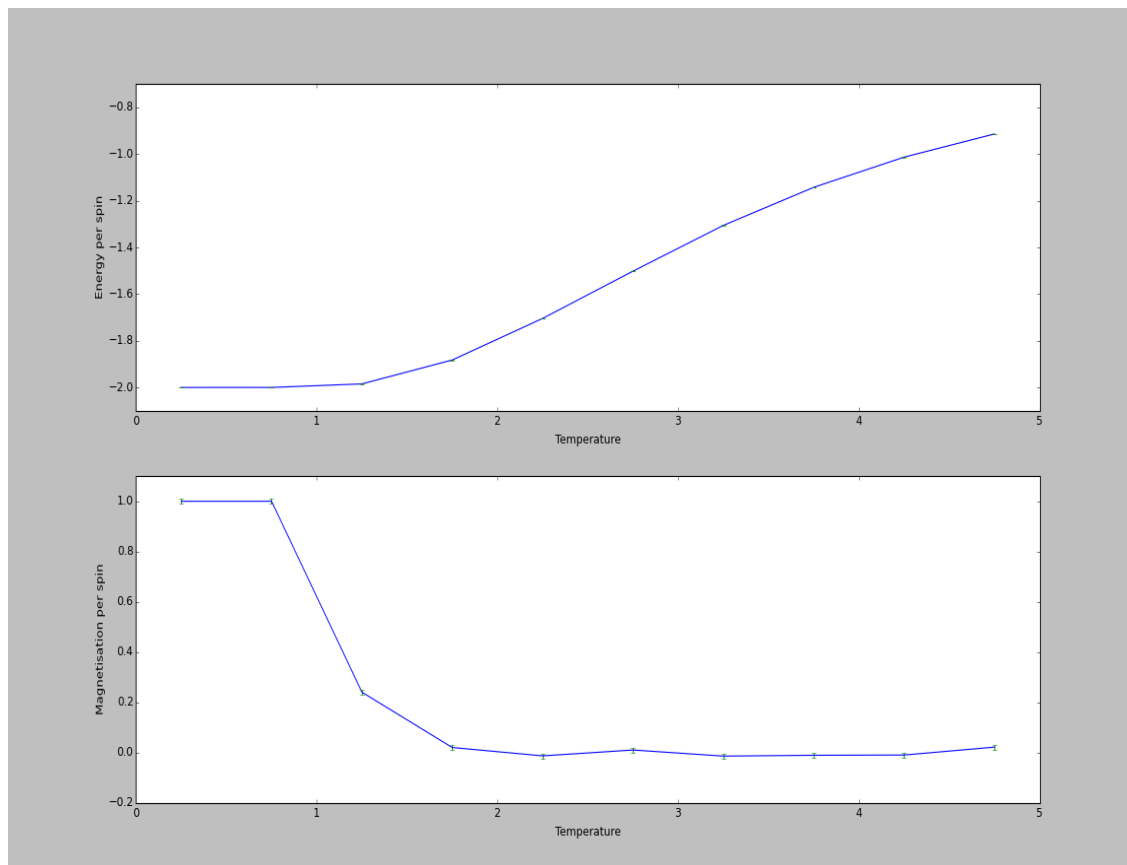


Figure 7. Plot showing Energy per spin vs Temperature for 8x8 lattice

Scaling the System Size

TASK:

1. Repeat the final task of the previous section for the following lattice sizes: 2x2, 4x4, 8x8, 16x16, 32x32. Make sure that you name each datafile that you produce after the corresponding lattice size!
2. Write a Python script to make a plot showing the energy *per spin* versus temperature for each of your lattice sizes. Hint: the NumPy loadtxt function is the reverse of the savetxt function, and can be used to read your previously saved files into the script.
3. Repeat this for the magnetisation. As before, use the plot controls to save your a PNG image of your plot and attach this to the report.
4. How big a lattice do you think is big enough to capture the long range fluctuations?

1.

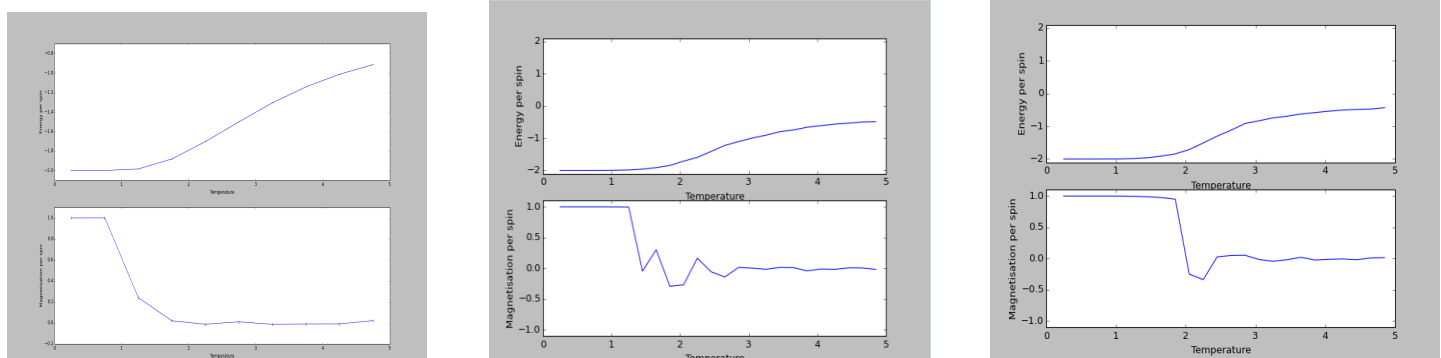


Figure 8. Examples of Energy per spin vs temperature for different lattice sizes. From left: 2x2, 4x4 8x8

It can be seen from figure 8 that the magnetisation is high before the curie temperature, and 0 after. The energy increases with temperature as expected.

2. Code can be found here:

[ComparingSystemSizes.py](#)

output from script:

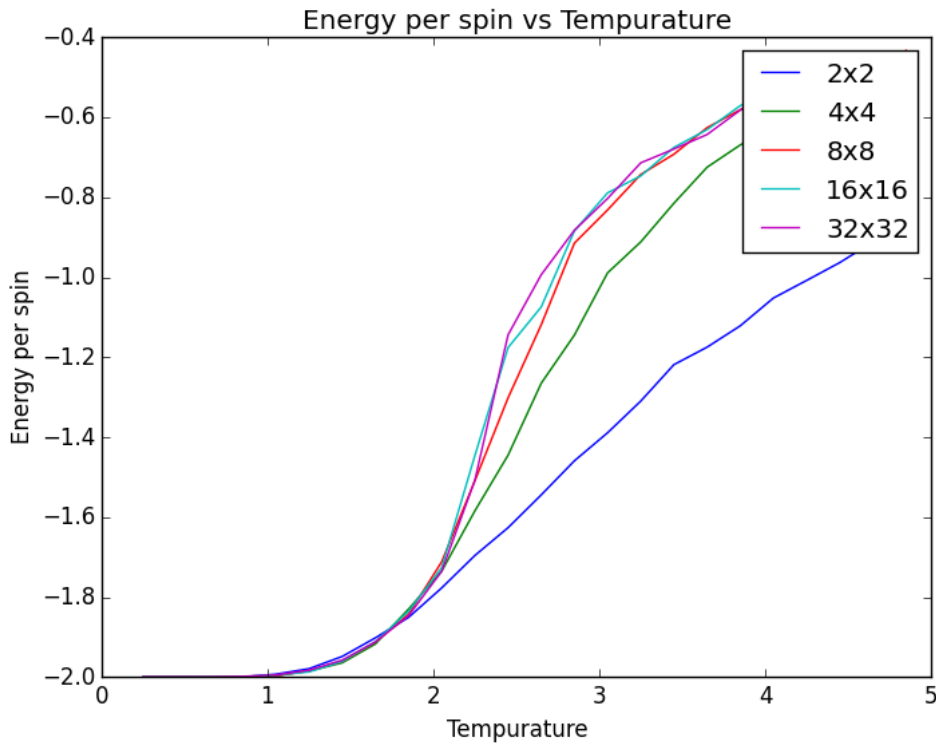


Figure 9. Energy per spin vs temperature for lattice sizes 2x2, 4x4, 8x8, 16x16 and 32x32

3.

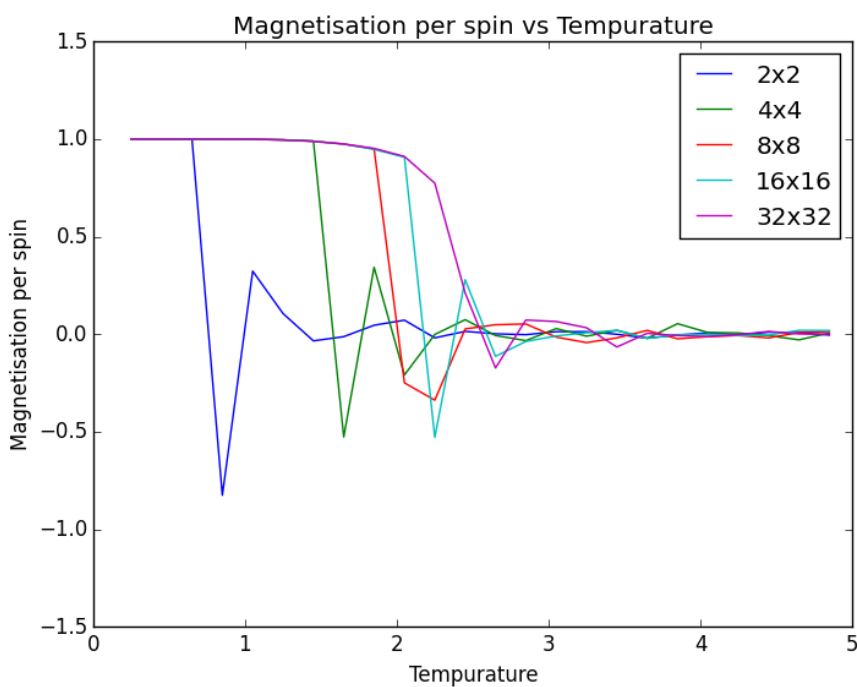


Figure 10. Magnetisation per spin vs temperature for lattice sizes 2x2, 4x4, 8x8, 16x16 and 32x32

4. The Energy per spin vs temperature figure shows that the gradient of the transition energy reaches an asymptote. Larger lattice sizes take longer time so 8x8 is the smallest lattice size to capture long range fluctuations.

Calculating the heat capacity

TASK: By definition, $C = \frac{\partial \langle E \rangle}{\partial T}$ From this, show that $C = \frac{\text{Var}[E]}{k_B T^2}$.

$$C = \frac{\partial \langle E \rangle}{\partial T}$$

$$\langle E \rangle = \frac{1}{Z} \sum_i E_i e^{-\frac{E_i}{k_B T}} \quad (1)$$

Z – Normalisation factor so when integrated over all space equals 1.

$$Z = \sum_i e^{-\frac{E_i}{k_B T}} \quad (2)$$

combining (1) and (2) gives:

$$\langle E \rangle = \frac{1}{\sum_i e^{-\frac{E_i}{k_B T}}} \sum_i E_i e^{-\frac{E_i}{k_B T}}$$

therefore:

$$C = \frac{\partial \langle E \rangle}{\partial T} = \frac{\partial}{\partial T} \left(\frac{\sum_i E_i e^{-\frac{E_i}{k_B T}}}{\sum_i e^{-\frac{E_i}{k_B T}}} \right)$$

Substituting

$$A = \sum_i E_i e^{-\frac{E_i}{k_B T}}$$

and

$$Z = \sum_i e^{-\frac{E_i}{k_B T}}$$

$$\frac{\partial \langle E \rangle}{\partial T} = \frac{\partial}{\partial T} \left(\frac{A}{Z} \right) = -\frac{A}{Z^2} \frac{\partial Z}{\partial T} + \frac{1}{Z} \frac{\partial A}{\partial T}$$

$$\frac{\partial Z}{\partial T} = \sum_i \frac{E_i}{k_B T^2} e^{-\frac{E_i}{k_B T}}$$

$$\frac{\partial A}{\partial T} = \sum_i \frac{E_i^2}{k_B T^2} e^{-\frac{E_i}{k_B T}}$$

Therefore, substituting gives:

$$\frac{\partial \langle E \rangle}{\partial T} = \frac{\partial}{\partial T} \left(\frac{A}{Z} \right) = -\frac{A}{Z^2} \sum_i \frac{E_i}{k_B T^2} e^{-\frac{E_i}{k_B T}} + \frac{1}{Z} \sum_i \frac{E_i^2}{k_B T^2} e^{-\frac{E_i}{k_B T}}$$

$$\begin{aligned}
&= \frac{1}{k_B T^2} \left(-\frac{A}{Z} \sum_i E_i e^{-\frac{E_i}{k_B T}} + \frac{1}{Z} \sum_i E_i^2 e^{-\frac{E_i}{k_B T}} \right) \\
&= \frac{1}{k_B T^2} \left(-\frac{A^2}{Z^2} + \sum_i \frac{1}{Z} E_i^2 e^{-\frac{E_i}{k_B T}} \right) \\
&= \frac{1}{k_B T^2} (-\langle E \rangle^2 + \langle E^2 \rangle)
\end{aligned}$$

Since:

$$\text{Var}[E] = \langle E^2 \rangle - \langle E \rangle^2$$

Gives the relationship:

$$\frac{\partial \langle E \rangle}{\partial T} = \frac{\text{Var}[E]}{k_B T}$$

TASK: Write a Python script to make a plot showing the heat capacity versus temperature for each of your lattice sizes from the previous section. You may need to do some research to recall the connection

between the variance of a variable, $\text{Var}[X]$, the mean of its square $\langle X^2 \rangle$, and its squared

mean $\langle X \rangle^2$. You may find that the data around the peak is very noisy — this is normal, and is a result of being in the critical region. As before, use the plot controls to save your a PNG image of your plot and attach this to the report.

The code can be found here:

[HeatCapacityVsTemp.py](#)

Output of code:

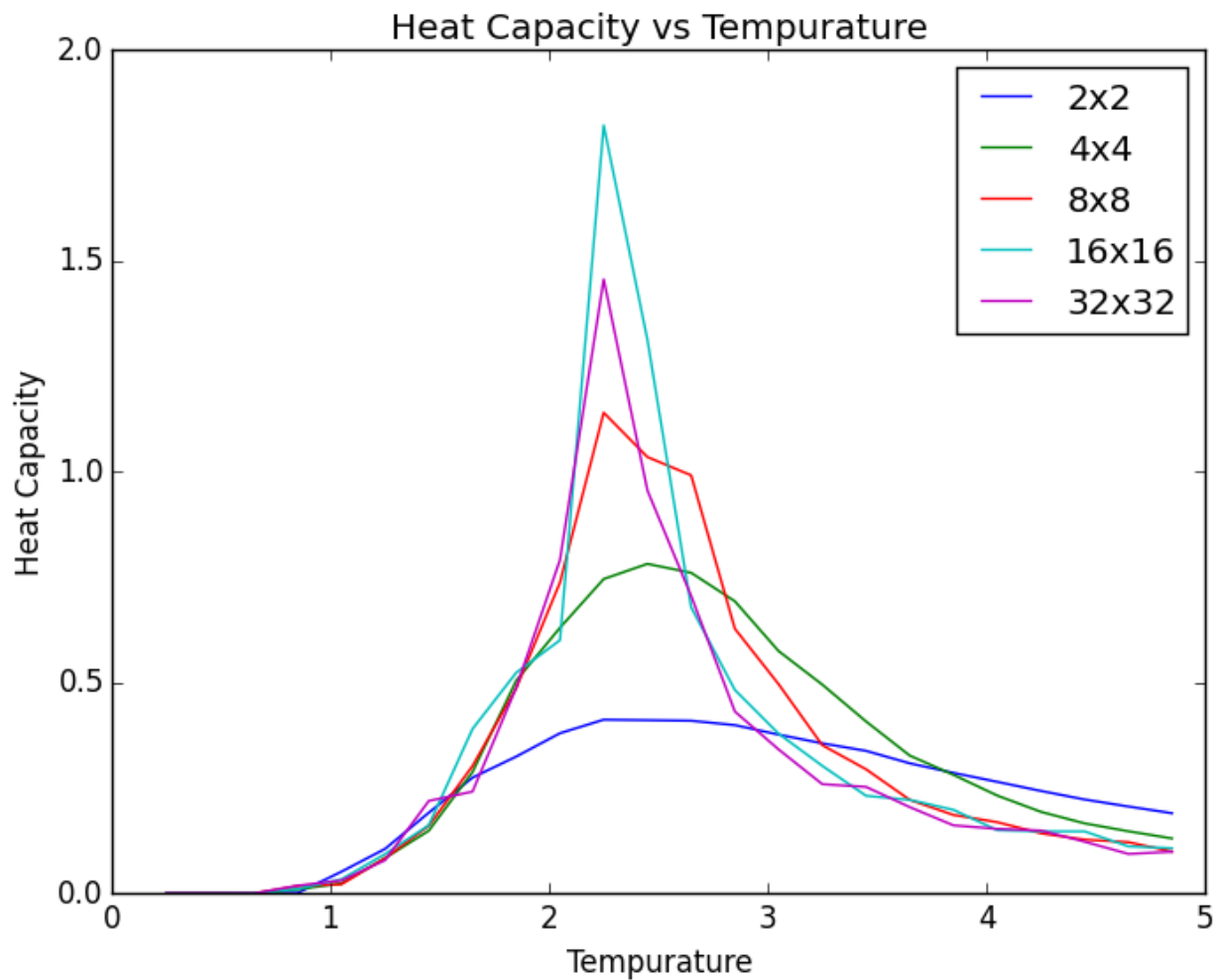


Figure 11. Heat capacity Vs Temperature for 2x2, 4x4, 8x8, 16x16, 32x32 lattice sizes

From figure 11 it can be seen that the heat capacity diverges as the lattice size increases, indicating that the heat capacity of an infinite 2D lattice can be computed.

TASK: A C++ program has been used to run some much longer simulations than would be possible on the college computers in Python. You can view its source code [here](#) if you are interested. Each file contains six columns: T , E , E^2 , M , M^2 , C (the final five quantities are per spin), and you can read them with the NumPy loadtxt function as before. For each lattice size, plot the C++ data against your data. For *one* lattice size, save a PNG of this comparison and add it to your report — add a legend to the graph to label which is which. To do this, you will need to pass the label="..." keyword to the plot function, then call the legend() function of the axis object (documentation [here](#)).

The code can be found here:

[CPlusPlusVsMyData.py](#)

the 2x2 lattice output is:

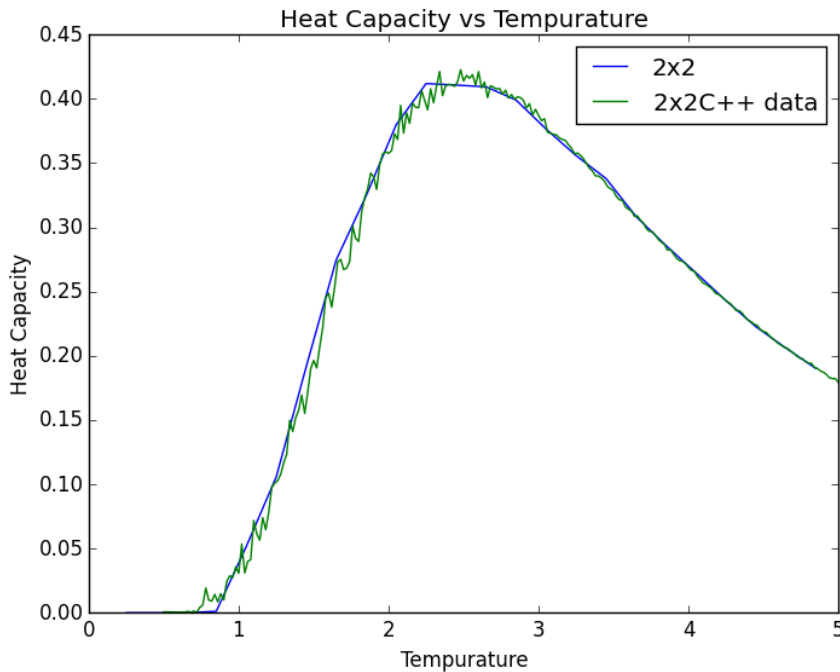


Figure 12. Heat capacity vs temperature for the 2x2 lattice. C++ data plotted against the python data.

TASK: write a script to read the data from a particular file, and plot C vs T, as well as a fitted polynomial. Try changing the degree of the polynomial to improve the fit — in general, it might be difficult to get a good fit! Attach a PNG of an example fit to your report.

The code can be found here:

[CPlusPlusVsMyDataPolyfit.py](#)

An example is shown below for the 2x2 lattice:

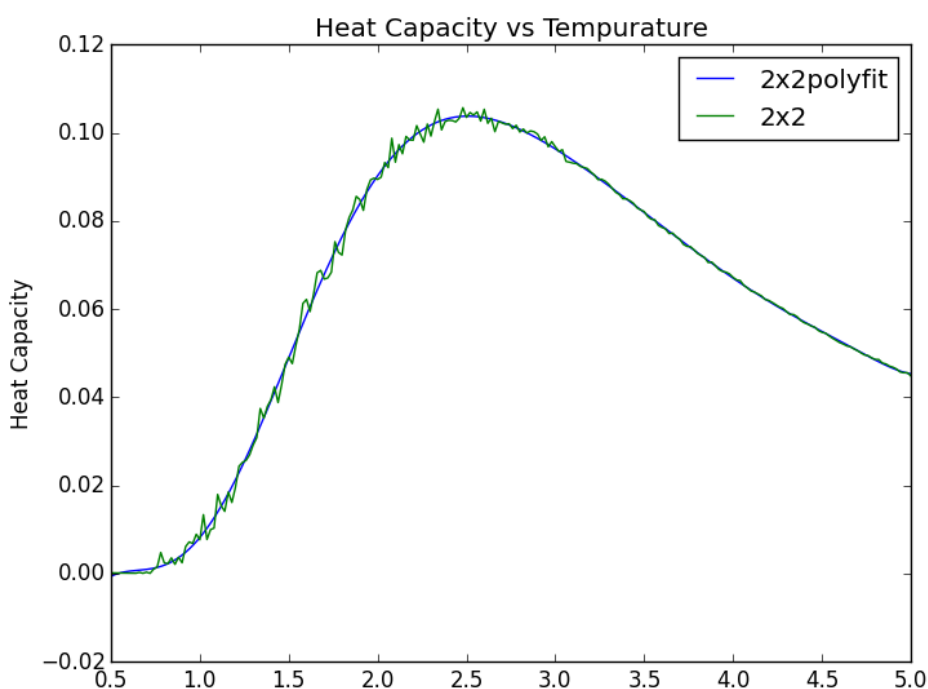


Figure 13. Heat capacity vs temperature for the 2x2 lattice using polyfit for the C++ data.

The degree of polynomial that fitted best was 10, however in the bigger lattices (16x16 and 32x32) the fit was not very accurate.

TASK: Modify your script from the previous section. You should still plot the whole temperature range, but fit the polynomial only to the peak of the heat capacity! You should find it easier to get a good fit when restricted to this region.

The code can be found here:

[CPlusPlusVsMyDataPolyfitPart2.py](#)

An example is shown below for the 2x2 lattice:

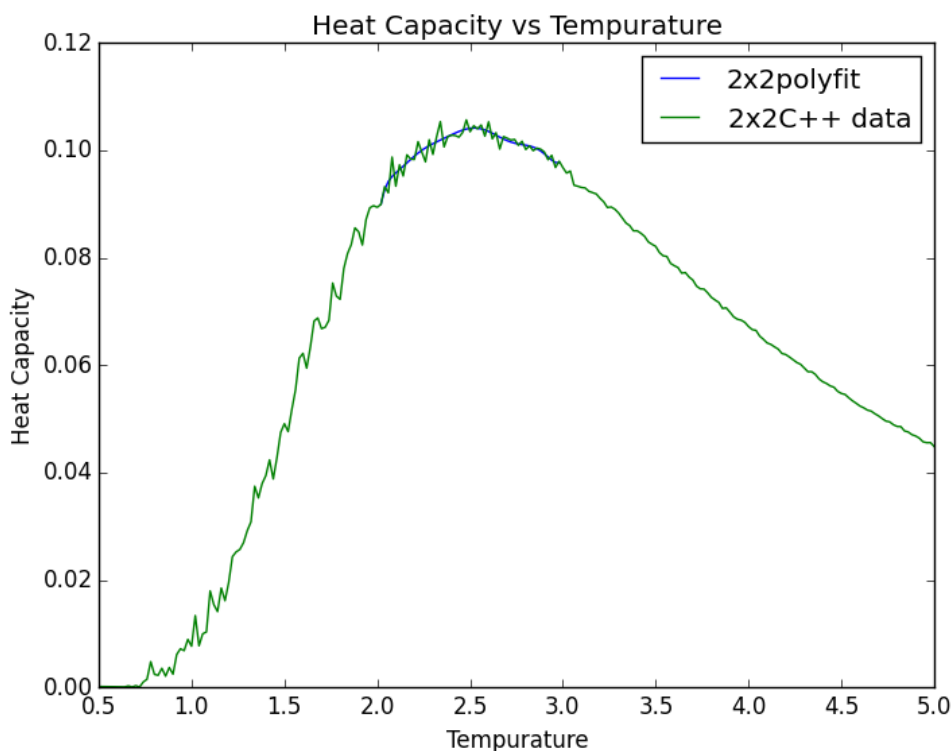


Figure 14. Heat capacity vs temperature for the 2x2 lattice using a restricted region for the polyfit.

There is a good fit.

TASK:

1. Find the temperature at which the maximum in C occurs for each datafile that you were given.
2. Make a text file containing two columns: the lattice side length (2,4,8, etc.), and the temperature at which C is a maximum. This is your estimate of T_C for that side length.
3. Make a plot that uses the scaling relation given above to determine $T_{C,\infty}$. By doing a little research online, you should be able to find the theoretical exact Curie temperature for the infinite 2D Ising lattice.
4. How does your value compare to this?
5. Are you surprised by how good/bad the agreement is?
6. Attach a PNG of this final graph to your report, and discuss briefly what you think the major sources of error are in your estimate.

1. The code can be found here (at the bottom):

[CPlusPlusVsMyDataPolyfitPart2.py](#)

Lattice Size	Max Temperature/K (T_C)
2x2	2.52
4x4	2.46
8x8	2.32
16x16	2.31
32x32	2.29

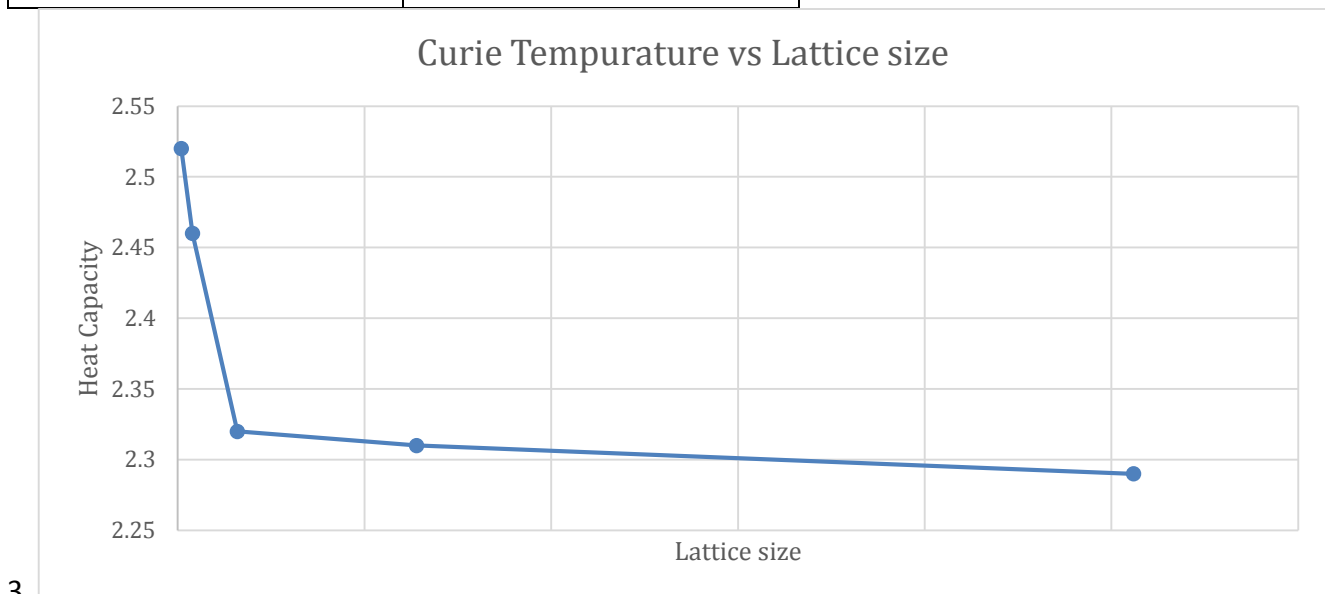


Figure 15. Graph showing trend in heat capacity against Lattice size

The graph reaches an asymptote at **~2.275**. This is determined by observation as no trend lines match this graph well.

4. Reference for duality argument (1):

$$\frac{k_B T_C}{J} = \frac{2}{\log(1 + \sqrt{2})} = 2.269 \dots$$

5. This is very similar to the Curie temperature giving an error of 2.275 ± 0.006 , which is much smaller than I was expecting!

6. The main sources of error in this experiment are:

1. The choice of 4000 as the cut off limit for the number of steps before taking the average: used 8x8 when 32x32 is likely to take more steps to reach equilibrium. A better alternative is for the cut off the change depending on lattice size.
2. Choice of degree of polynomial in polyfit to find max C
3. Trend line to find Heat capacity at infinity was by observation.

References:

- (1) 1941, H.A. Kramers and G.H. Wannier, *Phys. Rev.* **60**, 252 (1941)